# TODA FRAMES, HARMONIC MAPS AND EXTENDED DYNKIN DIAGRAMS

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ABSTRACT. We prove that all immersions of a genus one surface into G/T possessing a Toda frame can be constructed by integrating a pair of commuting vector fields on a finite dimensional Lie algebra. Here G is any simple real Lie group (not necessarily compact), T is a Cartan subgroup and the k-symmetric space structure on G/T is induced from the Coxeter automorphism. We provide necessary and sufficient conditions for the existence of a Toda frame for a harmonic map into G/T and describe those G/T to which the theory applies in terms of involutions of extended Dynkin diagrams.

# 1. INTRODUCTION

The last few decades have seen significant progress in the understanding and classification of harmonic maps from surfaces into compact real Lie groups and symmetric spaces. An important class of harmonic maps are those of *finite type*, which are obtained as the solutions to a pair of ordinary differential equations on a finite dimensional loop algebra. This is a far simpler process than attempting to solve the Laplace-Beltrami equation directly, and so motivates us to determine circumstances under which harmonic maps are of finite type. Similarly, when the target manifold is a k-symmetric space, k > 2, it is natural to restrict our attention to those harmonic maps which are cyclic primitive and ask when these maps are of finite type. Many papers (e.g. [10, 14, 2, 9, 4, 3, 6]) have addressed these questions when the target Lie group or (k)-symmetric space is *compact*. We remove the need for this compactness assumption and in Theorem 5.2 show that all maps from a genus one surface into a k-symmetric space G/T possessing a Toda frame are of finite type, where G is any simple real Lie group preserved by a Coxeter automorphism and T is the corresponding Cartan subgroup. A natural generalisation of the usual 2-dimensional affine Toda field equations provides the integrability condition for the existence of a Toda frame, and so we make contact with classical integrable systems theory. To determine the spaces G/T and the harmonic maps into them to which this theory applies we address the following two questions, each of independent interest:

- (1) When does a map from a surface into G/T possess a Toda frame? and
- (2) When is G preserved by a Coxeter automorphism?

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The first of these is answered in Theorem 4.2, where it is proven that a map from a surface into G/T locally has a Toda frame precisely when it is cyclic primitive and a certain function is constant. Cyclic primitive maps are in particular harmonic and play an analogous role for k-symmetric spaces as harmonic maps do for symmetric spaces. This and our finite-type result are the natural extensions of results obtained in [3] in the case when G is compact. The second question does not arise in the compact situation, since a Coxeter automorphism for a complex simple Lie algebra  $\mathfrak{g}^{\mathbb{C}}$  automatically preserves a compact real form  $\mathfrak{g}$ . We characterise when a Coxeter automorphism preserves a real form of a complex simple Lie algebra, which is equivalent to the corresponding real Lie group G being preserved whenever Gsimply connected or adjoint. Given simple roots for  $\mathfrak{g}^{\mathbb{C}}$  spanning a Cartan subalgebra  $\mathfrak{t}^{\mathbb{C}}$ , let  $\sigma$  be the associated Coxeter automorphism and  $\Theta$  a Cartan involution with respect to  $\mathfrak{g}$  that preserves  $\mathfrak{t} = \mathfrak{g} \cap \mathfrak{t}^{\mathbb{C}}$ . Then  $\sigma$  preserves  $\mathfrak{g}$  if and only if  $\Theta$ defines a permutation of the extended Dynkin diagram, so in particular whenever t is a maximally compact Cartan subalgebra (Proposition 3.1). In Theorem 3.2 we prove that all involutions of the extended Dynkin diagram for a simple complex Lie algebra  $\mathfrak{g}^{\mathbb{C}}$  arise from a Cartan involution for some real form  $\mathfrak{g}$ .

Harmonic maps from surfaces into Lie groups and symmetric spaces arise naturally in many geometric and physical problems. On the geometric side, strong motivation comes from the study of surfaces with particular curvature properties. For example, minimal surfaces are described by conformal harmonic maps and both constant mean curvature and Willmore surfaces are characterised by having harmonic Gauss maps into particular symmetric spaces. From the physics viewpoint, these harmonic maps are interesting because of their relationship with the appropriate Yang-Mills equations and non-linear sigma-models. Indeed the harmonic map equations on a Riemann surface are precisely the reduction of the Yang-Mills equations on  $\mathbb{R}^{2,2}$ obtained by considering solutions invariant under translation in the directions of negative signature. Classical solutions of sigma-models are given by harmonic maps into (non-compact) as pseudo-Riemannian manifolds. In [7] we study an explicit example, namely harmonic tori in de Sitter spaces  $S_1^m$ . In particular we apply the theory of this paper to the superconformal such maps with globally defined harmonic sequence to see that they may all be obtained by integrating a pair of commuting vector fields on a finite-dimensional vector space. It follows that all Willmore tori in  $S^3$  without umbilic points may be obtained in this simple way.

The structure of this paper is as follows. In section 2 we give the general theory for harmonic maps of surfaces into symmetric spaces and for primitive maps into k-symmetric spaces when the relevant Lie group G is equipped with a bi-invariant pseudo-metric. The question of when a Coxeter automorphism preserves the real form of the complex simple Lie algebra is addressed in section 3 in terms of Cartan involutions and extended Dynkin diagrams. Section 4 contains the relationship with the affine Toda field equations and the finite type result is proven in section 5.

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#### 2. FINITE TYPE MAPS INTO SYMMETRIC SPACES

The fact that a harmonic map from a surface to a Lie group corresponds to a loop of flat connections [15, 17] is the fundamental observation that enables one to apply integrable systems techniques to the study of these maps. The Cartan map  $G/H \to G$  from a symmetric space to the relevant Lie group is well-known to be a totally geodesic immersion when G is compact and equipped with a bi-invariant Riemannian metric. The composition of a harmonic map with a totally geodesic one is again harmonic, so this enables harmonic maps into symmetric spaces to be studied using the same tools as those into Lie groups, and in particular in terms of a loop of flat connections. We show in Theorem 2.1 that when G has merely a biinvariant pseudo-metric that the Cartan map is again a totally geodesic immersion. In particular all reductive Lie groups possess a bi-invariant pseudo-metric. We can hence study harmonic maps into G/H using integrable systems methods regardless of whether G is compact.

Let G be a semisimple Lie group. Recall that a homogeneous space G/H is a k-symmetric space (k > 1) if there is an automorphism  $\tau : G \to G$  of order k such that

$$(G^{\tau})_0 \subset H \subset G^{\tau}$$

where  $G^{\tau}$  denotes the fixed point set of  $\tau$ , and  $(G^{\tau})_0$  the identity component of  $G^{\tau}$ . When k = 2, we say that G/H is a symmetric space. We have the induced action

$$au: G/H \to G/H$$
 $gH \mapsto au(g)H.$ 

We write  $\tau$  also for the induced automorphism of  $\mathfrak{g}$  and note the  $\mathbb{Z}_k$ -grading

$$\mathfrak{g}^{\mathbb{C}} = igoplus_{j=0}^{k-1} \mathfrak{g}_j^{ au}, \; [\mathfrak{g}_j^{ au}, \mathfrak{g}_l^{ au}] \subset \mathfrak{g}_{j+l}^{ au},$$

where  $\mathfrak{g}_i^{\tau}$  denotes the  $e^{j\frac{2\pi i}{k}}$ -eigenspace of  $\tau$ .

We shall be interested in harmonic maps from a Riemann surface  $\Sigma$  into a symmetric space G/H. When G is compact, the Killing form on  $\mathfrak{g}$  induces a bi-invariant metric on G/H and the harmonic map equations for  $f: \Sigma \to G/H$  may either be calculated directly [20], using Noether's Theorem [16], or by composing f with the Cartan map  $G/H \to G$ , which is well-known in this case to be a totally geodesic immersion [8]. Recall here that the Cartan map of a symmetric space is given by

$$\iota: G/H \to G$$
$$gH \mapsto \tau(g)g^{-1}$$

We suppose merely that G has a bi-invariant pseudo-metric. Then analogous computations hold; in particular we can reduce the problem to studying harmonic maps into the Lie group G due to the following result. **Theorem 2.1.** Let G be a semisimple Lie group with bi-invariant pseudo-metric  $\langle \cdot, \cdot \rangle$  and G/H a symmetric space with respect to the involution  $\tau : G \to G$ . Then  $\iota : gH \mapsto \tau(g)g^{-1}$  is a totally geodesic immersion  $G/H \to G$ . If  $H = G^{\tau}$ , then  $\iota$  is additionally an embedding.

Let us call a Lie group G reductive if its Lie algebra  $\mathfrak{g}$  is reductive, that is has radical equal to its centre. Then  $\mathfrak{g}$  may be written as the direct sum of a semisimple Lie algebra and an abelian one. On the semisimple Lie algebra the Cartan-Killing form is non-degenerate, whilst on the abelian algebra any bilinear form is invariant under the adjoint action of the group. Combining these we obtain the existence of a bi-invariant pseudometric on any reductive Lie group, and hence the above theorem in particular applies when G is reductive.

Proof.  $\iota$  is an immersion: Suppose  $d\iota_{gH}(\gamma'(0)) = 0$  for some smooth path  $\gamma$  in G/H with  $\gamma(0) = gH$ . Take a lift  $\tilde{\gamma}$  of  $\gamma$  to G with  $\tilde{\gamma}(0) = g$  and write  $\pi : G \to G/H$  for the projection. Then

$$0 = \left. \frac{d}{dt} \right|_{t=0} \left( \tau\left(\tilde{\gamma}(t)\right) (\tilde{\gamma}(t))^{-1} \right) = d\tau_g(\tilde{\gamma}'(0))g^{-1} - \tau(g)g^{-1}\tilde{\gamma}'(0)g^{-1},$$

 $\mathbf{SO}$ 

$$d\tau_e(g^{-1}\tilde{\gamma}'(0)) = \tau(g^{-1})d\tau_g(\tilde{\gamma}'(0)) = g^{-1}\tilde{\gamma}'(0)$$

and  $\gamma'(0)$  is zero in  $T_{gH}(G/H)$  so  $d\iota_{gH}$  is injective.

 $\iota$  is totally geodesic: Let  $\nabla^l$  denote the connection on G obtained by trivialising TG by left translation, and similarly  $\nabla^r$  that induced from trivialising by right translation. A computation shows that  $\nabla^r = \nabla^l + \mathrm{ad}_{g^{-1}dg}$  and hence

$$\nabla = \frac{1}{2} (\nabla^l + \nabla^r)$$

is the Levi-Civita connection of the pseudo-metric  $\langle \cdot, \cdot \rangle$ .

Denote by exp :  $\mathfrak{g} \to G$  the Lie-theoretic exponential map, and by e the differentialgeometric exponential map associated to the Levi-Civita connection  $\nabla$ . Note that as in the definite case, for each  $X \in \mathfrak{g}$  the map

$$\gamma_X:\mathfrak{g}\to G$$
$$t\mapsto e^{tX}$$

is a geodesic, i.e.  $\nabla_{\gamma'_X} \gamma'_X = 0$ , so exp and e agree on the domain of e. Since the pseudo-metric is bi-invariant, we conclude that the geodesics through  $g \in G$  are locally of the form  $\gamma(t) = ge^{tX}$ . Denote by  $\mathfrak{m}$  the (-1)-eigenspace of  $\tau : \mathfrak{g} \to \mathfrak{g}$ , and note that  $\mathfrak{g} = \mathfrak{h} \oplus \mathfrak{m}$ , where  $\mathfrak{h}$  is the Lie algebra of H. The lift  $\tilde{\gamma}(t) = ge^{tX}H$  is horizontal, in the sense that  $\tilde{\gamma}'(t) \in ge^{tX}\mathfrak{m}$ . Thus the geodesics in G/H through gH are locally of the form  $\tilde{\gamma}(t) = ge^{tX}H$ . Since

$$\iota(ge^{tX}H) = ge^{tX}\tau(e^{-tX})\tau(g^{-1}) = ge^{2tX}\tau(g^{-1}) = g\tau(g^{-1})e^{t\tau(g)X\tau(g^{-1})}$$

is again a geodesic, we conclude that  $\iota$  is totally geodesic.

If  $H = G^{\tau}$ , then  $\iota$  is an embedding: In this case if  $\iota(g_1H) = \iota(g_2H)$ , then  $g_1^{-1}g_2 = \tau(g_1^{-1}g_2)$ , and so  $g_1^{-1}g_2 \in H$ , and thus  $\iota$  is injective.  $\Box$ 

Let  $F: U \to G$  be a smooth lift of  $f: U \to G/H$  on some simply connected  $U \subset \Sigma$ , where we assume henceforth that G is semisimple and has a bi-invariant pseudo-metric (we will later restrict our attention to simple such G.). By the above theorem, f is harmonic if and only if  $\iota \circ f$  is. The Maurer-Cartan form on G is the unique left-invariant  $\mathfrak{g}$ -valued 1-form which acts as the identity on  $\mathfrak{g}$ . We denote it by  $\omega$ , and note that if G is a linear group, then  $\omega = g^{-1}dg$ . We will use this notation throughout even in the non-linear case. Write  $\tilde{f} = \iota \circ f$  and  $\Phi = \tilde{f}^*(\omega) = \tilde{f}^{-1}d\tilde{f}$ . For any smooth  $\tilde{f}$ , the form  $\Phi$  satisfies the zero-curvature condition

(1) 
$$d\Phi + \frac{1}{2}[\Phi \wedge \Phi] = 0,$$

known as the Maurer-Cartan equation. Recall that for vector fields X, Y,

$$[\Phi \land \Phi](X,Y) = 2[\Phi,\Phi](X,Y) = [\Phi(X),\Phi(Y)].$$

The condition that the map  $\tilde{f}: \Sigma \to G$  is harmonic can be written as

(2) 
$$d * \Phi = 0.$$

Noting that  $\tilde{f} = \tau(F)F^{-1}$ , we have

(3) 
$$\Phi = F\left(\tau(F)^{-1}d(\tau(F)) - F^{-1}dF\right)F^{-1} = -2\mathrm{Ad}_F(\varphi_\mathfrak{m})$$

where  $\varphi = \varphi_{\mathfrak{h}} + \varphi_{\mathfrak{m}}$  is the decomposition of  $\varphi := F^{-1}dF$  into the eigenspaces of  $\tau$ . Then (2) becomes

(4) 
$$0 = d(\operatorname{Ad}_F(\ast\varphi_{\mathfrak{m}})) = \operatorname{Ad}_F(d \ast \varphi_{\mathfrak{m}} + [\varphi \land \ast\varphi_{\mathfrak{m}}])$$

or equivalently,

(5) 
$$d * \varphi_{\mathfrak{m}} + [\varphi \wedge * \varphi_{\mathfrak{m}}] = 0.$$

One can also compute the harmonic map equations directly for f. Writing  $[\mathfrak{m}]$  for the subbundle of  $G/H \times \mathfrak{g}$  whose fibre at  $g \cdot x$  is  $\operatorname{Ad}_g(\mathfrak{m})$ , we have an isomorphism  $[\mathfrak{m}] \cong T(G/H)$  given by

$$\begin{split} [\mathfrak{m}]_y &\to T_y G/H \\ Y &\mapsto \left. \frac{d}{dt} \right|_{t=0} e^{tY} \cdot y. \end{split}$$

The inverse of this isomorphism defines a  $\mathfrak{g}$ -valued 1-form  $\theta$  on the symmetric space G/H, which we term its Maurer-Cartan form. Then [16] f is harmonic if and only if

$$d*(f^*\theta)=0$$

and using that

$$f^*\theta = \mathrm{Ad}_F(\varphi_\mathfrak{m})$$

we recover (4). Write  $\varphi'_{\mathfrak{m}} + \varphi''_{\mathfrak{m}}$  for the decomposition of  $\varphi_{\mathfrak{m}}$  into dz and  $d\bar{z}$  parts. Since  $[\mathfrak{m}, \mathfrak{m}] \subset \mathfrak{h}$ , a straightforward computation shows (1) and (5) are equivalent to the requirement that for each  $\lambda \in S^1$ , the form

(6) 
$$\varphi_{\lambda} = \lambda \varphi'_{\mathfrak{m}} + \varphi_{\mathfrak{h}} + \lambda^{-1} \varphi''_{\mathfrak{m}}$$

satisfies the Maurer-Cartan equation

(7) 
$$d\varphi_{\lambda} + \frac{1}{2}[\varphi_{\lambda} \wedge \varphi_{\lambda}] = 0$$

Some solutions to (7) can be obtained simply by solving a pair of commuting ordinary differential equations on a finite-dimensional loop algebra. These unusually simple solutions are said to be of finite type.

Let G/K be a k-symmetric space for k > 2 and  $\tau$  the corresponding kth order involution. As we shall now explain when mapping into a k-symmetric space for k > 2 it is natural to restrict our attention to a subclass of harmonic maps consisting of those which are primitive, a notion that we now define. Again we have the reductive splitting

$$\mathfrak{g} = \mathfrak{k} \oplus \mathfrak{p}$$

with

$$\mathfrak{p}^{\mathbb{C}} = \bigoplus_{j=1}^{k-1} \mathfrak{g}_j^{\tau}, \qquad \mathfrak{k}^{\mathbb{C}} = \mathfrak{g}_0^{\tau}.$$

Similarly to before we may define the Maurer-Cartan form  $\theta$  of the k-symmetric space G/K when k > 2. For any smooth lift  $F: U \to G$  of  $\psi: U \to G/K$ , writing  $\varphi = F^* \omega$  we have

$$\psi^*\theta = \operatorname{Ad}_F \varphi_{\mathfrak{p}}$$

We say that a smooth map  $\psi$  of a surface  $\Sigma$  into G/K is *primitive* if the image of  $\psi^*\theta'$  is contained in  $[\mathfrak{g}_1]$ . Equivalently, it is primitive precisely when  $\varphi' = F^{-1}\partial F$  takes values in  $\mathfrak{g}_0^\tau \oplus \mathfrak{g}_1^\tau$ . Using that  $[\mathfrak{g}_1^\tau, \mathfrak{g}_{-1}^\tau] \subset \mathfrak{g}_0^\tau$ , the Maurer-Cartan equation for  $\varphi$  decomposes into  $\mathfrak{g}_1^\tau, \mathfrak{g}_0^\tau$  and  $\mathfrak{g}_{-1}^\tau$  components as

(8)  

$$d\varphi'_{\mathfrak{p}} + [\varphi_{\mathfrak{k}} \wedge \varphi'_{\mathfrak{p}}] = 0$$

$$d\varphi_{\mathfrak{k}} + \frac{1}{2} [\varphi_{\mathfrak{k}} \wedge \varphi_{\mathfrak{k}}] + [\varphi'_{\mathfrak{p}} \wedge \varphi''_{\mathfrak{p}}] = 0$$

$$d\varphi''_{\mathfrak{p}} + [\varphi_{\mathfrak{k}} \wedge \varphi''_{\mathfrak{p}}] = 0.$$

From these equations one easily verifies that primitive maps are in particular harmonic. Moreover [5] if G/H is a symmetric space with  $K \subset H$  and the corresponding reductive splitting preserved under  $\tau$ , then the projection of  $\psi : \Sigma \to G/K$  into G/His harmonic. An analogous calculation to that above shows that on simply connected subsets  $U \subset \Sigma$ , a primitive map  $\psi : U \to G/K$  is equivalent to a loop

(9) 
$$\varphi_{\lambda} = \lambda \varphi'_{\mathfrak{p}} + \varphi_{\mathfrak{k}} + \lambda^{-1} \varphi''_{\mathfrak{p}}, \quad \lambda \in S^{1}$$

of  $\mathfrak{g}$ -valued 1-forms each satisfying the Maurer-Cartan equation. We see then that both harmonic maps into symmetric spaces and primitive maps into k-symmetric

spaces are governed by the same equation (7) so we turn now to the question of constructing solutions to this equation.

Let  $\Omega G$  be the loop group  $\Omega G = \{\gamma : S^1 \to G\}$  with corresponding loop algebra  $\Omega \mathfrak{g} := \{\xi : S^1 \to \mathfrak{g}\}$ , where the loops are assumed real analytic without further comment. We use  $\Omega \mathfrak{g}^{\mathbb{C}}$  to denote loops in the complexified Lie algebra  $\mathfrak{g}^{\mathbb{C}}$ . For studying maps into k-symmetric spaces it is helpful to consider the twisted loop group

$$\Omega^{\tau}G = \{\gamma: S^1 \to G: \gamma(e^{\frac{2\pi i}{k}\lambda}) = \tau(\gamma(\lambda))\}$$

and corresponding twisted loop algebra  $\Omega^{\tau} \mathfrak{g}$  along with its complexification  $\Omega^{\tau} \mathfrak{g}^{\mathbb{C}}$ . The (possibly doubly infinite) Laurent expansion

$$\xi(\lambda) = \sum_{j} \xi_{j} \lambda^{j}, \quad \xi_{j} \in \mathfrak{g}_{j}^{\tau} \subset \mathfrak{g}^{\mathbb{C}}, \quad \Phi_{-j} = \bar{\Phi}_{j}$$

allows us to filtrate  $\Omega^{\tau}\mathfrak{g}^{\mathbb{C}}$  by finite-dimensional subspaces

$$\Omega_d^{\tau} = \{ \xi \in \Omega \mathfrak{g} \mid \xi_j = 0 \text{ whenever } |j| > d \}.$$

Fix a Cartan subalgebra  $\mathfrak{t}$  of  $\mathfrak{g}$  such that  $\mathfrak{t} \subset \mathfrak{k}$  and recall that a non-zero  $\alpha \in (\mathfrak{t}^{\mathbb{C}})^*$ is a root with corresponding root space  $\mathcal{G}^{\alpha} \subset \mathfrak{g}^{\mathbb{C}}$  if  $[X_1, X_2] = \alpha(X_1)X_2$  for all  $X_1 \in \mathfrak{t}$ and  $X_2 \in \mathcal{G}^{\alpha}$ . We denote the set of roots by  $\Delta$  and employ the same notation for the root system formed by considering  $\Delta$  as a subset of  $(\mathfrak{t}^{\mathbb{C}})^*$ . Choose a set of simple roots, that is a subset  $\{\alpha_1, \ldots, \alpha_N\}$  of  $\Delta$  such that every root  $\alpha \in \Delta$  can be written uniquely as

$$\alpha = \sum_{j=1}^{N} m_j \alpha_j,$$

where the  $m_j$  are either all positive integers or all negative integers. The height of  $\alpha$  is  $h(\alpha) = \sum_{j=1}^{N} m_j$  and the root(s) of maximal height are called highest root(s) whilst those of minimal height are termed lowest root(s).

We similarly define the root spaces of  $\mathfrak{k}^{\mathbb{C}}$ . Let  $\mathfrak{n}$  be the nilpotent algebra consisting of the positive root spaces of  $\mathfrak{k}^{\mathbb{C}}$  with respect to a choice of simple roots and consider the resulting Iwasawa decomposition

(10) 
$$\mathfrak{k}^{\mathbb{C}} = \mathfrak{n} \oplus \mathfrak{t}^{\mathbb{C}} \oplus \bar{\mathfrak{n}}$$

of  $\mathfrak{k}^{\mathbb{C}}$ . Then for  $\eta \in \mathfrak{k}^{\mathbb{C}}$  and a local coordinate z on  $\Sigma$ , decomposing according to (10) we have

$$(\eta dz)_{\mathfrak{h}} = r(\eta)dz + r(\eta)d\bar{z}$$

where  $r: \mathfrak{k}^{\mathbb{C}} \to \mathfrak{k}^{\mathbb{C}}$  is defined by

(11) 
$$r(\eta) = \eta_{\bar{\mathfrak{n}}} + \frac{1}{2}\eta_{\mathfrak{k}}.$$

The key observation here is that for simply-connected coordinate neighbourhood  $U \subset \Sigma$ , if  $\xi : U \to \Omega_d^{\tau}$  satisfies

(12) 
$$\frac{\partial \xi}{\partial z} = [\xi, \lambda \xi_d + r(\xi_{d-1})]$$

then

$$\varphi_{\lambda} = (\lambda \xi_d + r(\xi_{d-1}))dz + (\lambda^{-1}\xi_{-d} + \overline{r(\xi_{d-1})})d\overline{z}$$

satisfies the Maurer-Cartan equation (7) (c.f. [5], Theorem 2.5). The equation

$$\frac{1}{2}(X(\xi) - iY(\xi)) = (\lambda\xi_d + r(\xi_{d-1}))$$

defines vector fields X, Y on  $\Omega_d$ . A straightforward computation shows that these vector fields commute and so finding solutions to (12) is merely a matter of solving a pair of commuting ordinary differential equations. This yields a rather special class of solutions to the Maurer-Cartan equations (7) and hence of harmonic maps to symmetric spaces and primitive maps to k-symmetric spaces, k > 2. The flows of X, Y are easily seen to evolve on spheres in  $\Omega_d$ . When G is compact, so are these spheres and hence X, Y are complete and for any initial condition the differential equation (12) has a unique solution on U. However when G is non-compact the completeness of X, Y is not guaranteed.

**Definition 2.2.** A harmonic map  $f: \Sigma \to G/H$  to a symmetric space or a primitive map  $\psi: \Sigma \to G/K$  to a k-symmetric space, k > 2 is said to be of *finite type* if it has a lift  $F: \Sigma \to G$  for which there exists a smooth map  $\xi: \mathbb{R} \to \Omega_d^{\tau} \mathfrak{g}$  satisfying

(13) 
$$d\xi = [\xi, \varphi_{\lambda}]$$

and

(14) 
$$\varphi_{\lambda} = (\lambda \xi_d + r(\xi_{d-1}))dz + (\lambda^{-1}\xi_{-d} + \overline{r(\xi_{d-1})})d\overline{z}.$$

Here  $\varphi_{\lambda}$  and r are defined in (9) and (11) for the primitive case and in (6) and the obvious analogue to (11) for the harmonic case.

We introduce some terminology for later use. A formal Killing field for f or  $\psi$  is a smooth map  $\xi : \Sigma \to \Omega^{\tau} \mathfrak{g}$  satisfying the Lax equation (13). When  $\xi$  takes values in some  $\Omega_d$  it is termed a polynomial Killing field of degree d and when it additionally satisfies (14) it is an adapted polynomial Killing field.

When the automorphism  $\tau : \mathfrak{g}^{\mathbb{C}} \to \mathfrak{g}^{\mathbb{C}}$  is of the form  $\tau = \operatorname{Ad}_{\exp M}$  for some  $M \in \mathfrak{t}^{\mathbb{C}}$ where  $\mathfrak{t}$  is a Cartan subalgebra of  $\mathfrak{g}$ , then we can express the eigenspaces  $\mathfrak{g}_{j}^{\tau}$  of  $\tau$  in terms of root spaces.

Given our chosen set of simple roots  $\alpha_j$ , denote by  $\eta_j$  the corresponding dual basis of  $\mathfrak{t}^{\mathbb{C}}$ . For any root  $\alpha = m_1 \alpha_1 + \ldots m_N \alpha_N$ , smooth map  $s_j : \Sigma \to \mathbb{C}$  and root vector  $R_\alpha \in \mathcal{G}^\alpha$ , a straightforward computation shows that

(15) 
$$\operatorname{Ad}_{\exp(s_1\eta_1+\ldots s_N\eta_N)} R_{\alpha} = \exp(m_1s_1+\ldots m_Ns_N)R_{\alpha}.$$

Note that  $\exp(m_1 s_1 + \dots m_N s_N)$  is a scalar function. Given  $\tau = \operatorname{Ad}_{\exp(\frac{2\pi i}{k}(\sum s_j \eta_j))}$  we have

$$\mathfrak{g}_l^{\tau} = \operatorname{span}\{R_{\alpha} | \alpha = \sum_{j=1}^N m_j \alpha_j, \sum_{j=1}^N s_j m_j = l \bmod (k)\}.$$

In particular if we let k - 1 denote the maximal height of a root of  $\mathfrak{g}^{\mathbb{C}}$  and suppose (16)  $\sigma := \operatorname{Ad}_{\exp(\frac{2\pi i}{k} \sum_{j=1}^{N} \eta_j)},$ 

then  $\sigma$  is of order k and from (15) it acts on the root spaces by

(17) 
$$\sigma(R_{\alpha}) = \exp\left(\frac{2\pi i h(\alpha)}{k}\right) R_{\alpha}.$$

We recognise the inner automorphism  $\sigma$  as the Coxeter automorphism associated to the identity transformation of the simple roots [1]. It plays an important role here because when it preserves the real Lie group G, it allows us to view G/T as a k-symmetric space for which  $\mathfrak{g}_1^{\sigma}$  is the sum of the simple and lowest root spaces. Here T is a Cartan subgroup with Lie algebra  $\mathfrak{t}$ . Furthermore since K = T in this case, the map r described in (11) is simply multiplication by  $\frac{1}{2}$  and so the adapted polynomial Killing field condition (14) simplifies. Taking this N-symmetric space structure on G/T, we say that a smooth map  $\psi : \Sigma \to G/T$  is cyclic primitive if it is primitive and satisfies the condition that the image of  $\psi^*\theta'$  contains a cyclic element. Writing  $\alpha_0$  for the lowest root, an element in  $\left(\bigoplus_{j=0}^N \mathcal{G}^{\alpha_j}\right)$  is cyclic if its projection to each of the root spaces  $\mathcal{G}^{\alpha_0}, \mathcal{G}^{\alpha_1}, \ldots, \mathcal{G}^{\alpha_N}$  is non-zero. We henceforth assume that G is simple in order to guarantee the uniqueness of the lowest root (that is, we assume that G is connected and  $\mathfrak{g}$  is simple).

# 3. EXTENDED DYNKIN DIAGRAMS AND CARTAN INVOLUTIONS

We now ascertain the k-symmetric spaces to which our theory will apply. That is, we give conditions under which a choice of real form  $\mathfrak{g}$  of a simple complex Lie algebra  $\mathfrak{g}^{\mathbb{C}}$ , Cartan subalgebra  $\mathfrak{t}^{\mathbb{C}}$  and simple roots  $\alpha_j$  yield a Coxeter automorphism  $\sigma$  which preserves the real Lie algebra  $\mathfrak{g}$ . When  $G^{\mathbb{C}}$  is a simply connected or adjoint simple Lie group with Lie algebra  $\mathfrak{g}^{\mathbb{C}}$ , this ensures that the Coxeter automorphism preserves the real group G.

Let  $\bar{}$  denote the complex conjugation of  $\mathfrak{g}^{\mathbb{C}}$  corresponding to the real form  $\mathfrak{g}$ . Define the conjugate of a root  $\alpha$  by

$$\bar{\alpha}(X) = \alpha(\bar{X}).$$

Then from (17) we see that the condition for the Coxeter automorphism  $\sigma$  to preserve  $\mathfrak{g}$  is that for all roots  $\alpha$ , the height  $h(\alpha)$  satisfies

$$h(\bar{\alpha}) = -h(\alpha) \mod k$$

or equivalently that for  $j = 1, \ldots, N$  we have

$$\bar{\alpha}_j \in \{-\alpha_0,\ldots,-\alpha_N\}.$$

We will now use a Cartan involution to express this reality condition in terms of the extended Dynkin diagram for  $\alpha_0, \ldots, \alpha_N$ . A Cartan involution is an involution  $\Theta$  of  $\mathfrak{g}$  such that

$$\langle X, Y \rangle_{\Theta} = -\langle X, \Theta(Y) \rangle$$

is positive definite, where  $\langle \cdot, \cdot \rangle$  denotes the Killing form. Using complex-linearity,  $\Theta$  extends to an involution of  $\mathfrak{g}^{\mathbb{C}}$ . We may [12, Prop. 6.59] choose a Cartan involution  $\Theta$  which preserves the Cartan subalgebra  $\mathfrak{t}$ .

**Proposition 3.1.** Let  $\mathfrak{g}$  be a real simple Lie algebra,  $\mathfrak{t}$  a Cartan subalgebra and  $\Theta$  be a Cartan involution preserving  $\mathfrak{t}$ . Choose simple roots  $\alpha_1, \ldots, \alpha_N$  for the root system  $\Delta(\mathfrak{g}^{\mathbb{C}}, \mathfrak{t}^{\mathbb{C}})$  and let  $\sigma$  be the corresponding Coxeter automorphism of  $\mathfrak{g}^{\mathbb{C}}$  defined in (16). Then the following are equivalent:

- (1)  $\sigma$  preserves the real form  $\mathfrak{g}$ ,
- (2)  $\sigma$  commutes with  $\Theta$ ,
- (3)  $\Theta$  defines a permutation of the extended Dynkin diagram for  $\mathfrak{g}^{\mathbb{C}}$  consisting of the usual Dynkin diagram augmented with the lowest root  $\alpha_0$ .

*Proof.* Write  $\mathfrak{t} = \mathfrak{l} \oplus \mathfrak{p}$ , where  $\mathfrak{l}$ ,  $\mathfrak{p}$  are respectively the (+1)-eigenspace and (-1)eigenspace of the action of  $\Theta$  on  $\mathfrak{t}$ . Then [12, Cor. 6.49] all roots  $\alpha$  are real on  $\mathfrak{p}$ and imaginary on  $\mathfrak{l}$ , and defining the action of  $\Theta$  on roots by  $\Theta(\alpha)(X) = \alpha(\Theta(X))$ we have that

$$\Theta(\alpha) = -\bar{\alpha}$$
 for all roots  $\alpha$ .

If  $R_{\alpha}$  is a root vector for  $\alpha$ , then  $\overline{R}_{\alpha}$  is a root vector for  $\overline{\alpha}$  and  $\Theta(R_{\alpha})$  is a root vector for  $\Theta(\alpha)$ . We assume that our root vectors are chosen so that

$$R_{\bar{\alpha}} = R_{\alpha}$$

and write  $R_{\Theta(\alpha)} = c_{\alpha}\Theta(R_{\alpha})$ . Then using (17), a straightforward computation shows that  $\sigma \circ \Theta(R_{\alpha}) = \Theta \circ \sigma(R_{\alpha})$  if and only if  $\sigma(\bar{R}_{-\alpha}) = \overline{\sigma(R_{-\alpha})}$ , proving the equivalence of conditions (1) and (2) above.

The Cartan involution  $\Theta$  commutes with  $\sigma$  if and only if for all roots  $\alpha$ , the height function h satisfies

$$h(\Theta(\alpha)) \equiv h(\alpha) \mod k$$
,

or equivalently when  $\Theta$  defines a permutation of  $\alpha_0, \alpha_1, \ldots, \alpha_N$ . All automorphisms of a Lie algebra preserve the Killing form and hence a Cartan involution  $\Theta$  as above defines a permutation of the extended Dynkin diagram and we see the equivalence of conditions (2) and (3).

We next show that every involution of the extended Dynkin diagram for  $\Delta(\mathfrak{g}^{\mathbb{C}}, \mathfrak{t}^{\mathbb{C}})$ does indeed arise from a Cartan involution for some real form  $\mathfrak{g}$  with  $\Theta$ -stable Cartan subalgebra  $\mathfrak{t} = \mathfrak{g} \cap \mathfrak{t}^{\mathbb{C}}$ . A  $\Theta$ -stable Cartan subalgebra  $\mathfrak{t}$  of  $\mathfrak{g}$  is maximally compact if and only if  $\Theta$  preserves a set of simple roots for the root system  $\Delta(\mathfrak{g}, \mathfrak{t})$  [12, p 387]. Hence when  $\mathfrak{t}$  is maximally compact, a Coxeter automorphism  $\sigma$  must stabilise the real form  $\mathfrak{g}$ . (In particular, all Cartan subalgebras of a compact real form  $\mathfrak{g}$  are maximally compact.) The more interesting case then is when the Cartan subalgebra  $\mathfrak{t}$  is not maximally compact, which corresponds to the involution of the extended Dynkin diagram acting nontrivially on the lowest root  $\alpha_0$ .

**Theorem 3.2.** Every involution of the extended Dynkin diagram for a simple complex Lie algebra  $\mathfrak{g}^{\mathbb{C}}$  is induced by a Cartan involution of a real form of  $\mathfrak{g}^{\mathbb{C}}$ .

More precisely, let  $\mathfrak{g}^{\mathbb{C}}$  be a simple complex Lie algebra with Cartan subalgebra  $\mathfrak{t}^{\mathbb{C}}$  and choose simple roots  $\alpha_1, \ldots, \alpha_N$  for the root system  $\Delta(\mathfrak{g}^{\mathbb{C}}, \mathfrak{t}^{\mathbb{C}})$ . Given an involution  $\pi$ of the extended Dynkin diagram for  $\Delta$ , there exists a real form  $\mathfrak{g}$  of  $\mathfrak{g}^{\mathbb{C}}$  and a Cartan involution  $\Theta$  of  $\mathfrak{g}$  preserving  $\mathfrak{t} = \mathfrak{g} \cap \mathfrak{t}^{\mathbb{C}}$  such that  $\Theta$  induces  $\pi$  and  $\mathfrak{t}$  is a real form of  $\mathfrak{t}^{\mathbb{C}}$ . The Coxeter automorphism  $\sigma$  determined by  $\alpha_1, \ldots, \alpha_N$  preserves the real form  $\mathfrak{g}$ .

*Proof.* Let  $\pi$  be an involution of the extended Dynkin diagram. Denote also by  $\pi$  the corresponding involution of the set  $\{0, 1, \ldots, N\}$  and the induced involution of  $(\mathfrak{t}^{\mathbb{C}})^*$  which preserves the root system  $\Delta$  and satisfies  $\pi(\alpha_j) = \alpha_{\pi(j)}$ .

Let  $\{H_{\alpha}, R_{\alpha} \mid \alpha \in \Delta\}$  be a Chevalley basis. That is, writing  $\alpha^{\#}$  for the dual of the root  $\alpha$  with respect to the Killing form  $\kappa$  we set  $H_{\alpha} = (2/\kappa(\alpha^{\#}, \alpha^{\#}))\alpha^{\#}$  and we choose the root vectors  $R_{\alpha}$  so that

$$[R_{\alpha}, R_{-\alpha}] = H_{\alpha}.$$

and such that the structure constants  $c_{\alpha,\beta}$  defined by  $[R_{\alpha}, R_{\beta}] = c_{\alpha,\beta}R_{\alpha+\beta}$  satisfy  $c_{-\alpha,-\beta} = -c_{\alpha,\beta}$ . For any  $b_{\alpha_j} \in \mathbb{C}$  for  $j = 1, \ldots, N$ , we obtain an automorphism  $\Theta$  of  $\mathfrak{g}^{\mathbb{C}}$  compatible with  $\pi$  by requiring that  $\Theta(R_{\alpha_j}) = b_{\alpha_j}R_{\pi(\alpha_j)}$  for  $j = 1, \ldots, N$  and that  $\{\pi(H_{\alpha}), \Theta(R_{\alpha}) \mid \alpha \in \Delta\}$  is a Chevalley basis. Our first task is to verify that for an appropriate choice of  $b_{\alpha_j}$ , the resulting  $\Theta$  is involution.

Given  $\pi$  and  $b_{\alpha_1}, \ldots, b_{\alpha_N}$ , for any root  $\alpha$  we define  $b_{\alpha} \in \mathbb{C}$  by the equation  $\Theta(R_{\alpha}) = b_{\alpha}R_{\pi(\alpha)}$ . The automorphism  $\Theta$  will be an involution precisely when  $b_{\alpha_j}b_{\alpha_{\pi(j)}} = 1$  for  $j = 1, \ldots, N$ . For the j with  $\pi(j) \neq 0$ , we can clearly guarantee this by taking  $b_{\alpha_{\pi(j)}} = b_{\alpha_j}^{-1}$ . We will show that  $b_{\alpha_1}, \ldots, b_{\alpha_N}$  can be chosen so that additionally  $b_{\alpha_0}b_{\alpha_{\pi(0)}} = 1$ .

We may express  $R_{\alpha_0}$  as  $C[R_{-\beta_1}, [R_{-\beta_2}, \dots, [R_{-\beta_{K-1}}, R_{-\beta_K}] \dots]]$  for some non-zero constant C and simple roots  $\beta_i$  such that  $\sum_{i=1}^{K} \beta_i = -\alpha_0$ . Now writing  $\alpha_0 = -\sum_{j=1}^{N} m_j \alpha_j$  we have

(18) 
$$b_{\alpha_0} R_{\alpha_{\pi(0)}} = \Theta(R_{\alpha_0}) = C \prod_{j=1}^N b_{-\alpha_j}^{m_j} [R_{-\pi(\beta_1)}, [R_{-\pi(\beta_2)}, \dots, [R_{\pi(\beta_{K-1})}, R_{-\pi(\beta_K)}] \dots]]$$

and  $\Theta^2(R_{\alpha_0}) = \prod_{j=1}^N (b_{-\alpha_j} b_{-\alpha_{\pi(j)}})^{m_j} R_{\alpha_0}$ , implying

$$b_{\alpha_0}b_{\alpha_{\pi(0)}} = \prod_{j=1}^{N} (b_{-\alpha_j}b_{-\alpha_{\pi(j)}})^{m_j}.$$

Using that  $\{\pi(H_{\alpha}), \Theta(R_{\alpha}) \mid \alpha \in \Delta\}$  is again a Chevalley basis and that an automorphism of the extended Dynkin diagram must preserve the Killing form gives

$$b_{\alpha_j}b_{-\alpha_j} = \frac{\kappa(\pi(\alpha_j)), \pi(\alpha_j))}{\kappa(\alpha_j, \alpha_j)} = 1.$$

Hence

$$b_{\alpha_0}b_{\alpha_{\pi(0)}} = \prod_{j=1}^N (b_{\alpha_j}b_{\alpha_{\pi(j)}})^{-m_j} = (b_{\alpha_{\pi(0)}}b_{\alpha_0})^{-1},$$

where the last equality uses the assumption  $b_{\alpha_j}b_{\alpha_{\pi(j)}} = 1$  for  $\pi(j) \neq 0$ . We therefore automatically have  $b_{\alpha_0}b_{\alpha_{\pi(0)}} = \pm 1$ . Considering (18) shows that if there exists jsuch that  $\pi(j) = j$  and  $m_j$  is odd then by switching the sign of  $b_{\alpha_j}$  if necessary we may ensure that  $b_{\alpha_0}b_{\alpha_{\pi(0)}} = 1$ .

It remains to give a method of proof for when there is no  $\alpha_j$  with  $m_j$  odd that is fixed by  $\pi$ . If  $\pi(0) = 0$  then there is nothing to prove so we assume henceforth that  $\pi(0) \neq 0$ . Suppose  $\gamma$  is a positive root such that

(a) the expression  $\gamma = \sum_{j=1}^{N} n_j \alpha_j$  as a sum of simple roots has  $n_{\pi(0)} = 0$ , (b)  $\pi(\gamma) + \alpha_{\pi(0)}$  is also a root, and (c)  $\gamma + \alpha_0 = -\pi(\gamma) - \alpha_{\pi(0)}$ .

From (c) we have that

$$[[R_{\gamma}, R_{\alpha_0}], [R_{\pi(\gamma)}, R_{\pi(0)}]] = c_{\gamma, \alpha_0} c_{\pi(\gamma), \pi(\alpha_0)} H_{\gamma + \alpha_0}$$

Applying  $\Theta$  gives

$$\left[\left[b_{\gamma}R_{\pi(\gamma)}, b_{\alpha_0}R_{\pi(0)}\right], \left[b_{\pi(\gamma)}R_{\gamma}, b_{\alpha_{\pi(0)}}R_{\alpha_0}\right]\right] = -c_{\gamma,\alpha_0}c_{\pi(\gamma),\pi(\alpha_0)}H_{\gamma+\alpha_0}$$

and so

(19) 
$$b_{\gamma}b_{\pi(\gamma)}b_{\alpha_0}b_{\alpha_{\pi(0)}} = 1.$$

We may write  $R_{\gamma}$  as  $C'[R_{\beta'_1}, [R_{\beta'_2} \dots [R_{\beta'_{K'-1}}, R_{\beta'_{K'}}]] \dots]$  with C' a non-zero constant and  $\beta'_i \neq \alpha_{\pi(0)}$  simple roots satisfying  $\sum_{i=1}^{K'} \beta'_i = \gamma$ . Then

$$b_{\gamma}b_{\pi(\gamma)}R_{\gamma} = \Theta^2(R_{\gamma}) = \left(\prod_{i=1}^{K'} b_{\beta'_i}b_{\beta'_{\pi(i)}}\right)R_{\gamma}$$

However for simple roots  $\alpha_j$  with  $\pi(j) \neq 0$  we chose  $b_{\alpha_j}$  so that  $b_{\alpha_j} b_{\alpha_{\pi(j)}} = 1$  and hence  $b_{\gamma} b_{\pi(\gamma)} = 1$ . Substituting this into (19) gives that  $b_{\alpha_0} b_{\alpha_{\pi(0)}} = 1$ , as required. A similar argument applies if there are positive roots  $\gamma, \delta$  such that

- (i) the expressions of  $\gamma, \delta$  as sums of simple roots do not contain  $\alpha_{\pi(0)}$ ,
- (ii)  $\pi(\gamma) + \alpha_{\pi(0)}$  and  $\delta + \pi(\delta)$  are also roots, and
- (iii)  $\delta + \pi(\delta) + \gamma + \pi(\gamma) = -\alpha_0 \alpha_{\pi(0)}$ .

Here we know there is some non-zero constant C'' such that

$$[[R_{\gamma}, R_{\alpha_0}], [R_{\pi(\gamma)}, R_{\pi(0)}]] = C''[R_{-\delta}, R_{-\pi(\delta)}]$$

and as above applying  $\Theta$  gives

$$b_{\gamma}b_{\pi(\gamma)}b_{\alpha_0}b_{\alpha_{\pi(0)}} = b_{-\delta}b_{-\pi(\delta)}.$$

By (i) we know  $b_{\gamma}b_{\pi(\gamma)} = 1$  and  $b_{-\delta}b_{-\pi(\delta)} = 1$  so conclude that  $b_{\alpha_{\pi(0)}}b_{\alpha_0} = 1$ .

To show that every involution of the extended Dynkin diagram extends to an involution of the Lie algebra we now consider the involutions of each of the diagrams and, for those that do not fix some  $\alpha_j$  with odd  $m_j$ , identify a suitable root  $\gamma$  or pair of roots  $\gamma, \delta$ .

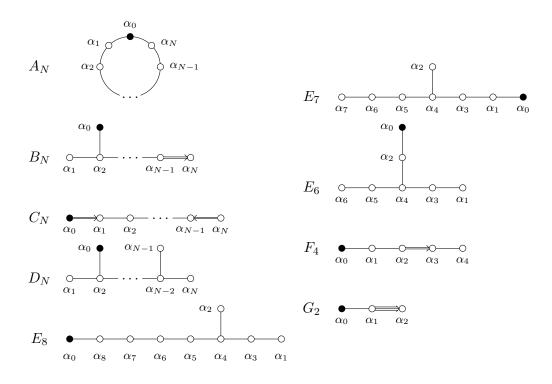


FIGURE 1. Extended Dynkin diagrams, with the lowest root  $\alpha_0$  coloured.

For a root system of type  $A_N$ , the simple root coefficients  $m_j = 1$  for all j = 1, ..., N. Thus any diagram involution fixing some node is induced by an involution of the Lie algebra. By inspection of the extended Dynkin diagram shown in Figure 3, we see that when N is even, every involution of the extended Dynkin diagram fixes some  $\alpha_j$ . When N is odd we need to consider the rotation  $\pi(j) = j + \frac{1}{2}(N+1) \mod (N+1)$ and reflections.

For the involution  $\pi(j) = j + \frac{1}{2}(N+1) \mod (N+1)$ , the root  $\gamma = \alpha_1 + \alpha_2 + \ldots + \alpha_{\frac{1}{2}(N-1)}$  satisfies conditions (a), (b), (c) above.

Consider now an involution  $\pi$  coming from a reflection. Since we have automatically covered the cases when there is a fixed root we can assume that there are an even number of roots between  $\alpha_0$  and  $\pi(\alpha_0)$  going in each direction around the circle. Indeed the axis of reflection is between the nodes  $(\pi(0) - 1)/2$  and  $(\pi(0) + 1)/2$  and between  $(N + \pi(0))/2$  and  $(N + \pi(0))/2 + 1$ . The roots

$$\gamma = \alpha_1 + \alpha_2 + \ldots + \alpha_{(\pi(0)-1)/2}$$
 and  $\delta = \alpha_{\pi(0)+1} + \alpha_{\pi(0)+2} \dots + \alpha_{(\pi(0)+N)/2}$ 

satisfy conditions (i), (ii), (iii) above.

There is only one involution of the root system of type  $B_N$ , which sends  $\alpha_0$  to  $\alpha_1$ and fixes everything else. We can choose  $\gamma = \alpha_2 + \ldots + \alpha_N$ .

For root systems of type  $C_N$  there is again only one involution;  $\pi(\alpha_i) = \alpha_{N-i}$ . Here choose  $\gamma = \alpha_1 + \ldots + \alpha_{N-1}$ .

For  $D_N$ ,  $m_1 = m_{N-1} = m_N = 1$ , and so we need only consider involutions which do not fix any of these vertices, of which there are three. These are involutions with  $\pi(0) = 1, N-1$  or N. If  $\pi(0) = 1$  then let  $\gamma = \alpha_2 + \ldots + \alpha_{N-1}$ , and if  $\pi(0) = N-1$ or N, take  $\gamma = \alpha_1 + \alpha_2 + \ldots + \alpha_{N-2}$ .

For the root system  $E_6$ , all involutions of the diagram fix the vertex  $\alpha_4$  and  $m_4 = 3$  is odd.

The unique involution of the extended Dynkin diagram for  $E_7$  satisfies  $\pi(\alpha_0) = \alpha_7$ . A list of all positive roots of  $E_7$  are tabulated for example in [18, p 1524-1530]. Let  $\gamma = \alpha_1 + \alpha_2 + 2\alpha_3 + 2\alpha_4 + \alpha_5 + \alpha_6$ , so  $\pi(\gamma) = \alpha_1 + \alpha_2 + \alpha_3 + 2\alpha_4 + 2\alpha_5 + \alpha_6$  and  $\pi(\gamma) + \alpha_{\pi(0)} = \alpha_1 + \alpha_2 + \alpha_3 + 2\alpha_4 + 2\alpha_5 + \alpha_6 + \alpha_7$  is also a root. Furthermore  $\gamma + \alpha_{\pi(0)} + \pi(\gamma) = 2\alpha_1 + 2\alpha_2 + 3\alpha_3 + 4\alpha_4 + 3\alpha_5 + 2\alpha_6 + \alpha_7$  which is the highest root.

The extended Dynkin diagrams of type  $E_8, F_4, G_2$  do not possess any involutions.

We have then shown that given any involution  $\pi$  of an extended Dynkin diagram for  $(\mathfrak{g}^{\mathbb{C}},\mathfrak{t}^{\mathbb{C}})$ , there exists an involution  $\Theta$  of  $\mathfrak{g}^{\mathbb{C}}$  preserving  $\mathfrak{t}^{\mathbb{C}}$  and inducing  $\pi$ . It remains to show that there is a real form  $\mathfrak{g}$  of  $\mathfrak{g}^{\mathbb{C}}$  for which  $\Theta$  is a Cartan involution and such that  $\mathfrak{g} \cap \mathfrak{t}^{\mathbb{C}}$  has full rank. For any choice of simple roots we may consider the corresponding *Borel subalgebra*  $\mathfrak{b}^{\mathbb{C}} = \mathfrak{t}^{\mathbb{C}} \oplus \bigoplus_{\alpha \in \Delta^+} \mathcal{G}^{\alpha}$  and it is easy to see that  $\Theta$  preserves the set of simple roots if and only if it preserves the corresponding Cartan and Borel subalgebras. Now by [11, Theorem 8.6] there exists an automorphism  $\Psi$  of  $\mathfrak{g}^{\mathbb{C}}$  such that  $\Psi \Theta \Psi^{-1}$  acts on the corresponding simple and lowest root vectors  $R_{\alpha_j}$  in the Chevalley basis simply by scaling them by  $\pm 1$ , and hence preserves the Cartan subalgebras  $\Psi^{-1}(\mathfrak{t}^{\mathbb{C}})$  and the Borel subalgebras  $\Psi^{-1}(\mathfrak{b}^{\mathbb{C}})$  and hence the set of simple roots are a real form  $\mathfrak{g}'$  of  $\mathfrak{g}^{\mathbb{C}}$  with respect to which  $\Theta$  is a Cartan involution and such that  $\mathfrak{g}'\Psi^{-1}(\mathfrak{t}^{\mathbb{C}})$  is a Cartan subalgebra of  $\mathfrak{g}'$  [12, proof of Theorem 6.88].

Let  $\mathfrak{l}^{\mathbb{C}}$  denote the (+1)-eigenspace of  $\Theta$  and  $L^{\mathbb{C}}$  a complex Lie group with Lie algebra  $\mathfrak{l}^{\mathbb{C}}$ . In [13, Theorem 1] (c.f. [19, Proposition 2.1]) it was shown that for a given real form  $\mathfrak{g}'$  and  $\Theta$ -stable Cartan subalgebra  $\mathfrak{t}^{\mathbb{C}}$  of a simple complex Lie algebra  $\mathfrak{g}^{\mathbb{C}}$ , there

exists a  $\Theta$ -stable Cartan subalgebra  $\mathfrak{t}'$  of  $\mathfrak{g}'$  and  $l \in L^{\mathbb{C}}$  such that  $\mathfrak{t}^{\mathbb{C}} = \mathrm{Ad}_{l}(\mathfrak{t}')^{\mathbb{C}}$ . Hence  $\mathfrak{g} = \mathrm{Ad}_{l} \mathfrak{g}'$  is a real form of  $\mathfrak{g}^{\mathbb{C}}$  for which  $\mathfrak{t} = \mathfrak{g} \cap \mathfrak{t}^{\mathbb{C}}$  is a  $\Theta$ -stable real form of  $\mathfrak{t}^{\mathbb{C}}$  and  $\Theta$  is a Cartan involution of  $\mathfrak{g}$ .

By Proposition 3.1 the Coxeter automorphism corresponding to the choice of simple roots  $\alpha_1, \ldots, \alpha_N$  preserves the real form  $\mathfrak{g}$  and in particular the Cartan subalgebra  $\mathfrak{t}$ .

#### 4. TODA FRAME

We now explore the relationship between cyclic primitive maps and the affine Toda field equations. Henceforth G shall denote a simple real Lie group, T a Cartan subgroup and  $\alpha_1, \ldots, \alpha_N$  simple roots such that the resulting Coxeter automorphism  $\sigma$  preserves the real group G. This Coxeter automorphism then gives G/T the structure of a k-symmetric space, where k-1 is the maximum height of a root of  $\mathfrak{g}^{\mathbb{C}}$ . We shall consider cyclic primitive maps  $\psi$  from the complex plane into G/T and will see that cyclic primitive maps  $\psi : \mathbb{C} \to G/T$  arise from and give rise to solutions of the two-dimensional affine Toda field equations for  $\mathfrak{g}$ . Our results also apply to maps from a simply-connected coordinate neighbourhood of any Riemann surface.

The famous Toda equations arose originally as a model for particle interactions within a one-dimensional crystal, with the affine model corresponding to the particles being arranged in a circle. They have been the subject of extensive study, both as a completely integrable Hamiltonian system and in the context of Toda field theories. The standard form of the affine Toda field equation for  $\mathfrak{g}$  on the complex plane is

(20) 
$$2\Omega_{z\bar{z}} = \sum_{j=0}^{N} m_j e^{2\alpha_j(\Omega)} \alpha_j^{\sharp}$$

Here  $\Omega : \mathbb{C} \to i\mathfrak{t}$  is a smooth map, the lowest root  $\alpha_0$  is given by

$$\alpha_0 = -\sum_{j=1}^N m_j \alpha_j,$$

we set  $m_0 = 1$  and  $R_{\alpha_j}$  are root vectors such that  $\alpha_j^{\sharp}$  is the dual of  $\alpha_j$  with respect to the Killing form. Using (3), since the Coxeter automorphism preserves the real form  $\mathfrak{g}$  there exists a permutation  $\pi$  of the roots  $\alpha_0, \alpha_1, \ldots, \alpha_N$  such that

(21) 
$$\overline{\alpha_j} = -\alpha_{\pi(j)}.$$

We shall consider the generalisation of the affine Toda field equations obtained by allowing  $m_0, m_1, \ldots, m_N$  to be any positive real numbers such that  $m_{\pi(j)} = \overline{m_j}$  and  $R_{\alpha_j}$  to be any root vectors satisfying  $\overline{R_{\alpha_j}} = R_{-\alpha_{\pi(j)}}$ .

Given a cyclic element  $W = \sum_{j=0}^{N} r_j R_{\alpha_j}$  of  $\mathfrak{g}_1^{\sigma}$ , we say that a lift  $F : \mathbb{C} \to G$  of  $\psi : \mathbb{C} \to G/T$  is a *Toda frame* with respect to W if there exists a smooth map

 $\Omega: \mathbb{C} \to i\mathfrak{t}$  such that

(22) 
$$F^{-1}F_z = \Omega_z + \operatorname{Ad}_{\exp\Omega} W$$

We call  $\Omega$  an *affine Toda field* with respect to W. The motivation for this nomenclature is

**Lemma 4.1.** Fix a cyclic element  $W = \sum_{j=0}^{N} r_j R_{\alpha_j}$  of  $\mathfrak{g}_1^{\sigma}$  such that  $m_{\pi(j)} = \overline{m_j}$ and  $\overline{R_{\alpha_j}} = R_{-\alpha_{\pi(j)}}$ .

The affine Toda field equation (20) is the integrability condition for the existence of a Toda frame with respect to W where we take  $m_j = r_j \overline{r_j}$  for j = 0, ..., N.

*Proof.* Using  $[R_{\alpha_j}, R_{-\alpha_l}] = 0$  whenever  $j \neq l$ , we can rewrite the Toda field equation (20) as

$$2\Omega_{z\bar{z}} = \sum_{j,l=0}^{N} r_j \overline{r_l} e^{\alpha_j(\Omega)} e^{\alpha_j(\Omega)} [R_{\alpha_j}, R_{-\alpha_l}]$$
$$= \left[\sum_{j=0}^{N} r_j e^{\alpha_j(\Omega)} R_{\alpha_j}, \sum_{l=0}^{N} \overline{r_l} e^{\alpha_l(\Omega)} R_{-\alpha_l}\right]$$

From equation (15) we know  $e^{\alpha_j(\Omega)}R_{\alpha_j} = \operatorname{Ad}_{\exp\Omega}R_{\alpha_j}$  and also

$$e^{\alpha_l(\Omega)}R_{-\alpha_l} = e^{-\alpha_l(-\Omega)}R_{-\alpha_l} = \operatorname{Ad}_{\exp-\Omega}R_{-\alpha_l}.$$

If we set  $W := \sum_{j=0}^{N} r_j R_{\alpha_j}$  with the normalisation is described in the lemma then since  $\sum_{j=0}^{N} R_j R_{\alpha_j} = \sum_{j=0}^{N} R_j \overline{R_{-\alpha_j}}$ , the Toda field equation becomes

$$2\Omega_{z\bar{z}} = [\operatorname{Ad}_{\exp\Omega} W, \operatorname{Ad}_{\exp-\Omega} \overline{W}].$$

Now for any given  $\Omega : \mathbb{C} \to i\mathfrak{t}$  the integrability condition for the existence of a Toda frame with respect to W is the Maurer-Cartan equation (1) for

$$\varphi = (\Omega_z + \operatorname{Ad}_{\exp\Omega} W)dz + (\Omega_{\bar{z}} + \operatorname{Ad}_{\exp-\Omega} W)d\bar{z}$$

Namely, this integrability condition is

$$0 = (-\Omega_{\bar{z}} + \operatorname{Ad}_{\exp -(\Omega)}\overline{W})_{z} - (\Omega_{z} + \operatorname{Ad}_{\exp \Omega}W)_{\bar{z}} + [\Omega_{z} + \operatorname{Ad}_{\exp \Omega}W, -\Omega_{\bar{z}} + \operatorname{Ad}_{\exp -(\Omega)}\overline{W}] = -2\Omega_{z\bar{z}} + [\operatorname{Ad}_{\exp \Omega}W, \operatorname{Ad}_{\exp -\Omega}\overline{W}],$$

which is precisely the Toda field equation.

Recall that we write  $\alpha_0 = -\sum_{j=1}^N m_j \alpha_j$  for the expression of the lowest root  $\alpha_0$  in terms of the chosen simple roots  $\alpha_1, \ldots, \alpha_N$ . Given  $\tilde{F} : \mathbb{C} \to G$  with

(23) 
$$\tilde{F}^{-1}\tilde{F}_z|_{\mathfrak{g}_1^{\sigma}} = \sum_{j=0}^{N} c_j R_{\alpha_j},$$

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we say that a cyclic element

(24) 
$$W = \sum_{j=0}^{N} r_j R_{\alpha_j}$$

 $\mathfrak{g}_1^{\sigma}$  is normalised with respect to  $\tilde{F}: \mathbb{C} \to G$  if

$$r_0 \prod_{j=1}^{N} r_j^{m_j} = c_0 \prod_{j=1}^{N} c_j^{m_j}$$

**Theorem 4.2.** A map  $\psi : \mathbb{C} \to G/T$  possesses a Toda frame if and only if it is cyclic primitive Let  $\psi : \mathbb{C} \to G/T$  be a cyclic primitive map possessing a frame  $\tilde{F} : \mathbb{C} \to G$ such that  $c_0 \prod_{j=1}^N c_j^{m_j}$  is constant, where  $c_j$  are the root coefficients defined in (23). Then for any cyclic element W of  $\mathfrak{g}_1^{\sigma}$  which is normalised with respect to  $\tilde{F}$  there exists a Toda frame  $F : \mathbb{C} \to G$  of  $\psi$  with respect to W. Furthermore if  $\psi$  and  $\tilde{F}$  are doubly periodic with lattice  $\Lambda$  then so is the Toda frame F.

Conversely, if  $\psi : \mathbb{C} \to G/T$  has a Toda frame F with respect to cyclic  $W \in \mathfrak{g}_1^{\sigma}$  then  $\psi$  is cyclic primitive and W is normalised with respect to F. In particular then the root coefficients  $c_j$  are such that  $c_0 \prod_{j=1}^N c_j^{m_j}$  is constant.

*Proof.* Consider the frames  $F := \tilde{F} \exp X$  of  $\psi$  where  $X : \mathbb{C} \to \mathfrak{t}$ . For such F we have  $F^{-1}F_z = \operatorname{Ad}_{\exp -X} \tilde{F}^{-1}\tilde{F}_z + X_z$  and so

$$F^{-1}F_z|_{\mathfrak{g}_1^{\sigma}} = \operatorname{Ad}_{\exp -X} \tilde{F}^{-1}\tilde{F}_z|_{\mathfrak{g}_1^{\sigma}}.$$

This implies the Toda condition of  $\operatorname{Ad}_{\exp\Omega} W = F^{-1}F_z|_{\mathfrak{g}_1^{\sigma}}$  is equivalent to

(25) 
$$\operatorname{Ad}_{\exp(X+\Omega)} W = \tilde{F}^{-1} \tilde{F}_z|_{\mathfrak{g}_1^{\sigma}} = \sum_{j=0}^N c_j R_{\alpha_j}.$$

Using equation (15) we can rewrite this as

$$\sum_{j=0}^{N} r_j e^{\alpha_j (X+\Omega)} R_{\alpha_j} = \sum_{j=0}^{N} c_j R_{\alpha_j}.$$

Comparing root space coefficients implies that

(26) 
$$e^{\alpha_j(X+\Omega)} = \frac{c_j}{r_j} \text{ for } j = 1, \dots k$$

and  $r_0 \prod_{j=1}^{N} (e^{\alpha_j (X+\Omega)})^{-m_j} = c_0$ . Since W is normalised with respect to  $\tilde{F}$  and  $\mathbb{C}$  is simply connected, we can solve for  $X + \Omega$ . We can then find  $\Omega$  and X from  $X + \Omega$  by taking its  $\mathfrak{t}$  and  $i\mathfrak{t}$  components respectively.

It remains to show that  $\Omega_z dz = F^{-1} \partial F|_{\mathfrak{t}} = \varphi'_{\mathfrak{t}}$ . From the  $\mathfrak{g}_1^{\sigma}$  component (8) of the Maurer-Cartan equation for  $\varphi$  we have

$$\partial(\operatorname{Ad}_{\exp\Omega} W) - [\operatorname{Ad}_{\exp\Omega} W, \varphi'_{\mathfrak{t}}] = 0$$

or equivalently

$$[\operatorname{Ad}_{\exp\Omega} W, \varphi'_{\mathfrak{t}} - \partial\Omega] = 0.$$

Since W is cyclic so is  $\operatorname{Ad}_{\exp\Omega} W$  and thus  $\varphi'_{\mathfrak{t}} = \partial\Omega$ .

Conversely, given W and a solution  $\Omega$  to the corresponding affine Toda field equation, the resulting Toda frame F is primitive. Furthermore the equation

$$r_0(e^{-\sum_{j=1}^N m_j \alpha_j (X+\Omega)} R_{\alpha_0} + \sum_{j=1}^N r_j e^{\alpha_j (X+\Omega)} R_{\alpha_j} = \sum_{j=0}^N c_j R_{\alpha_j}$$

implies that  $r_0 \prod_{j=1}^N r_j^{m_j} = c_0 \prod_{j=1}^N c_j^{m_j}$  and hence  $c_0 \prod_{j=1}^N c_j^{m_j}$  is a non-zero constant. This implies that the  $c_j$  are nowhere zero and  $\psi$  is cyclic primitive.

Now suppose  $\tilde{F}$  is doubly periodic with respect to a lattice  $\Lambda$ . Then for  $j = 1, \ldots N$ , from (26) we see that  $e^{\alpha_j(X+\Omega)}$  is doubly periodic with respect to  $\Lambda$  and so

$$\exp(X + \Omega) = \exp(\sum_{j=1}^{N} \alpha_j (X + \Omega) \eta_j)$$

is also. Given any  $\Gamma \in \Lambda$  it follows that

(27) 
$$\exp(X(z+\Gamma) - X(z)) = \exp(\Omega(z) - \Omega(z+\Gamma)).$$

Using the conjugation map  $\mathfrak{g}^{\mathbb{C}} \to \mathfrak{g}^{\mathbb{C}}$  which fixes  $\mathfrak{g}$ , we obtain from (27) that

(28) 
$$\exp(X(z+\Gamma) - X(z)) = \exp(-\Omega(z) + \Omega(z+\Gamma))$$

When combined, (27) and (28) imply that  $\exp(X(z + \Gamma)) = \exp(z))$  for all z and hence  $\exp X$  is doubly periodic with lattice  $\Lambda$ .

Since  $\tilde{F}$  and  $\exp X$  are both doubly periodic with lattice  $\Lambda$  we know  $F = \tilde{F} \exp X$  is also.

Our chief interest lies in cyclic primitive  $\psi$  which are doubly periodic, as it is these we shall show are of finite type. We henceforth restrict our attention to doubly-periodic maps and denote by  $\mathbb{C}/\Lambda$  any genus one Riemann surface. Let W be a cyclic element of  $\mathfrak{g}_1^{\sigma}$  as before. We say that a frame  $F : \mathbb{C}/\Lambda \to G$  of  $\psi : \mathbb{C}/\Lambda \to G/T$  is a Toda frame with respect to W if F is a Toda frame of  $\psi$  when both are considered as maps from  $\mathbb{C}$ . From the proof of Theorem 4.2 we make the following observation, which will prove useful in the next section.

**Lemma 4.3.** If  $F : \mathbb{C}/\Lambda \to G$  is a Toda frame of  $\psi : \mathbb{C}/\Lambda \to G/T$  then the corresponding affine Toda field  $\Omega : \mathbb{C} \to i\mathfrak{t}$  has the property that  $\exp \Omega$  and  $\Omega_z$  are doubly periodic with lattice  $\Lambda$ .

# 5. FINITE TYPE RESULT

We will now show that all smooth maps  $\psi$  from a 2-torus  $\mathbb{C}/\Lambda$  into the k-symmetric space G/T which have a Toda frame are of finite type. Hence all such maps can

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be constructed from a pair of commuting ordinary differential equations on a finitedimensional loop algebra. In [4] it was shown that all semisimple adapted harmonic maps of a 2-torus into a compact semisimple Lie group are of finite type. We prove our finite type result by adapting the methods of that paper. Note that the existence of a Toda frame forces  $\psi$  to be cyclic primitive.

A map  $Y : \mathbb{C}/\Lambda \to \mathfrak{g}^{\mathbb{C}}$  is called a *Jacobi field* if there exists  $\dot{\Omega} : \mathbb{C}/\Lambda \to \mathfrak{t}^{\mathbb{C}}$  such that

(29) 
$$dY + [F^{-1}dF, Y] = \left(\dot{\Omega}_z + [\dot{\Omega}, F^{-1}F_z]\right)dz + \left(-\dot{\Omega}_{\bar{z}} - [\dot{\Omega}, F^{-1}F_{\bar{z}}]\right)d\bar{z}.$$

If  $F_t$  is a family of Toda frames with corresponding  $\Omega_t : \mathbb{C} \to i\mathfrak{t}$  then  $\frac{d}{dt}F_t|_{t=0}$  is a Jacobi field with  $\dot{\Omega} = \frac{d}{dt}\Omega_t|_{t=0}$ . Note that if  $\dot{\Omega} = 0$  the Jacobi equation is the Killing field equation.

Let F be a Toda frame for  $\psi : \mathbb{C} \to G/T$ . We have

$$F^{-1}dF = (\Omega_z + \operatorname{Ad}_{\exp\Omega} W)dz + (-\Omega_{\overline{z}} + \operatorname{Ad}_{\exp-\Omega} \overline{W})d\overline{z}$$

for some  $\Omega: T^2 \to i\mathfrak{t}$  and cyclic  $W \in \mathfrak{g}_1^{\sigma}$ . Let Y be a Jacobi field with corresponding  $\dot{\Omega}: T^2 \to i\mathfrak{t}$ . Then Y must satisfy

(30) 
$$Y_z + [\Omega_z + \operatorname{Ad}_{\exp\Omega} W, Y] = \dot{\Omega}_z + [\dot{\Omega}, \operatorname{Ad}_{\exp\Omega} W]$$

(31) 
$$Y_{\bar{z}} + \left[-\Omega_{\bar{z}} + \operatorname{Ad}_{\exp-\Omega}\overline{W}, Y\right] = -\dot{\Omega}_{\bar{z}} - \left[\dot{\Omega}, \operatorname{Ad}_{\exp-\Omega}\overline{W}\right].$$

Taking  $(30)_{\bar{z}} - (31)_z$  we obtain

$$2\dot{\Omega}_{z\bar{z}} = -\left[\operatorname{Ad}_{\exp\Omega}W, [\dot{\Omega}, \operatorname{Ad}_{\exp-\Omega}\overline{W}]\right] - \left[\operatorname{Ad}_{\exp-\Omega}\overline{W}, [\dot{\Omega}, \operatorname{Ad}_{\exp\Omega}W]\right].$$

Since  $\Omega$  and W are fixed, we see that  $\dot{\Omega}$  satisfies a linear elliptic partial differential equation. As the torus is compact, the space of possible  $\dot{\Omega}$  is finite dimensional.

**Lemma 5.1.** Suppose  $\psi : \mathbb{C}/\Lambda \to G/T$  is a cyclic primitive map possessing a formal Killing field  $Y = \sum_{j \leq 1} \lambda^j Y_j \in \Omega^{\sigma} \mathfrak{g}^{\mathbb{C}}$ . Then  $\psi$  has a (real) polynomial Killing field with highest term  $Y_1$ .

*Proof.* We will find an infinite number of linearly independent Jacobi fields for which some linear combination must be a formal Killing field. Since Y is a formal Killing field, we have (13).

$$\sum_{j \leq 1} \lambda^j dY_j = \left\lfloor \sum_{j \leq 1} \lambda^j Y_j, \varphi_\lambda \right\rfloor.$$

Comparing coefficients of  $\lambda^j$  gives the equations

$$\begin{split} (Y_j)_z dz + [\varphi'_{\mathfrak{t}}, Y_j] + [\varphi'_{\mathfrak{p}}, Y_{j-1}] &= 0, \\ (Y_j)_{\bar{z}} d\bar{z} + [\varphi''_{\mathfrak{t}}, Y_j] + [\varphi''_{\mathfrak{p}}, Y_{j+1}] &= 0. \end{split}$$

For each  $l \in \mathbb{Z}^+$  set

$$Y^{l} := \frac{1}{2} Y_{-kl} + \sum_{-kl < j \le 1} \lambda^{j+kl} Y_{j}.$$

We will show that the  $Y^l$  are all Jacobi fields. Considering the coefficients separately gives

$$(Y^l)_z dz + [\lambda \varphi'_{\mathfrak{p}} + \varphi'_{\mathfrak{t}}, Y^l] = \frac{1}{2} (Y_{-kl})_z dz + \left\lfloor \frac{1}{2} Y_{-kl}, \lambda \varphi'_{\mathfrak{p}} \right\rfloor$$
$$(Y^l)_{\bar{z}} d\bar{z} + \left[ \varphi''_{\mathfrak{t}} + \lambda^{-1} \varphi''_{\mathfrak{p}}, Y^l \right] = -\frac{1}{2} (Y_{-kl})_{\bar{z}} d\bar{z} - \left\lfloor \frac{1}{2} Y_{-kl}, \lambda^{-1} \varphi''_{\mathfrak{p}} \right\rfloor.$$

Since  $Y_{-kl} \in \mathfrak{g}_0 = \mathfrak{t}^{\mathbb{C}}$  we can set  $\dot{\Omega}^l := \frac{1}{2}Y_{-kl}$ . With this choice of  $\dot{\Omega}$ , Y is a solution to (29) and hence is a Jacobi field. The space of potential  $\dot{\Omega}$  is finite dimensional, so there must be a non-trivial finite linear combination of the  $\dot{\Omega}^l$  which equals 0. The corresponding finite linear combination of the  $Y^l$  is a formal Killing field. Since the highest order terms of the  $Y^l$  are each  $Y_1$  we can rescale this formal Killing field to one with highest order term  $Y_1$ . After multiplying by an appropriate power of  $\lambda^k$  we may also assume that the degree of the lowest term has smaller absolute value than the degree of the highest term. Then  $\overline{\xi} + \xi$  is a polynomial Killing field for  $\xi$  and by construction has highest order term  $Y_1$ .

**Theorem 5.2.** Suppose  $\psi : \mathbb{C}/\Lambda \to G/T$  has a Toda frame  $F : \mathbb{C}/\Lambda \to G$ . Then  $\psi$  is of finite type.

*Proof.* Let  $F : \mathbb{C}/\Lambda \to G$  be the Toda frame of  $\psi$  with corresponding  $\Omega : \mathbb{C}/\Lambda \to i\mathfrak{t}$ and  $W \in \mathfrak{g}_1^{\sigma}$ . Recall that  $\psi$  is of finite type if it has an adapted polynomial Killing field  $\xi$ , that is a  $\xi = \sum_{j=-d}^d \lambda^j \xi_j$  in the real twisted loop algebra  $\Omega^{\sigma}\mathfrak{g}$  satisfying the Killing field equation (13) and such that

$$\xi_d + \lambda \frac{1}{2} \operatorname{Ad}_{\exp \Omega} W.$$

Since G was assumed simple, the complexified Lie algebra  $\mathfrak{g}^{\mathbb{C}}$  is simple and hence has a faithful linear representation so can be regarded as a subalgebra of some  $\mathfrak{gl}(m, \mathbb{C})$ . If we set

$$D = d - \mathrm{ad}_{\Omega_z dz - \Omega_{\bar{z}} d\bar{z}}$$

then we can rewrite (13) as

$$D\xi_{\lambda} = [\xi_{\lambda}, (2\Omega_z + \lambda \operatorname{Ad}_{\exp\Omega} W)dz + (-2\Omega_{\bar{z}} + \lambda^{-1}\operatorname{Ad}_{\exp-\Omega} \overline{W})d\bar{z}].$$

From  $d(\operatorname{Ad}_{\exp\Omega} W) = [\Omega_z, \operatorname{Ad}_{\exp\Omega} W]dz + [\operatorname{Ad}_{\exp\Omega} W, \Omega_{\bar{z}}]d\bar{z}$  we know  $D \operatorname{Ad}_{\exp\Omega} W = 0$ .

Writing  $V = \ker \operatorname{ad}_{\operatorname{Ad}_{\exp\Omega} W}$  and  $V^{\perp} = \operatorname{im} \operatorname{ad}_{\operatorname{Ad}_{\exp\Omega} W}$ , we have a bundle decomposition  $\mathbb{C}/\Lambda \times \mathfrak{g}^{\mathbb{C}} = V \oplus V^{\perp}$ . Furthermore

$$VV \subset V, \quad V^{\perp}V \subset V^{\perp}, \quad VV^{\perp} \subset V^{\perp},$$

Let  $X = \sum_{k \le -1} \lambda^k X_k$  where the  $X_k$  are sections of  $V^{\perp}$ . We seek X such that (32)  $Y = (1+X)^{-1} \operatorname{Ad}_{\exp \Omega} W(1+X)$  is a solution of the Killing field equation. Note that

$$DY = (1+X)^{-1} [\operatorname{Ad}_{\exp\Omega} W, DX(1+X)^{-1}](1+X),$$

and define a one-form  $\kappa$  by

$$\begin{split} \kappa &= (1+X)((2\Omega_z + \lambda\operatorname{Ad}_{\exp\Omega}W)dz + (-2\Omega_{\bar{z}} + \lambda^{-1}\operatorname{Ad}_{\exp-\Omega}\overline{W})d\bar{z} - DX)(1+X)^{-1}. \end{split}$$
 Routine calculations show that

$$DY + [(2\Omega_z + \lambda \operatorname{Ad}_{\exp\Omega} W)dz + (-2\Omega_{\bar{z}} + \lambda^{-1} \operatorname{Ad}_{\exp-\Omega} \overline{W})d\bar{z}, Y]$$
  
=  $(1 + X)^{-1}[\operatorname{Ad}_{\exp\Omega} W, -\kappa](1 + X)$ 

and hence Y satisfies the Killing field equation if and only if  $\kappa$  takes values in V. Our task then is to construct X so that  $\kappa$  takes values in V. We have

$$\kappa' \cdot (1+X) = (1+X)(2\Omega_z + \lambda \operatorname{Ad}_{\exp\Omega} W)dz - \partial X$$

where  $\kappa' \cdot (1+X)$  denotes multiplication. Note that  $\Omega_z$  is valued in  $V^{\perp}$  as it lies in  $\mathfrak{t}^{\mathbb{C}}$ .

The splitting of  $\kappa' \cdot (1+X)$  into its V and  $V^{\perp}$  components is

$$(V): \quad \kappa' = \left(\lambda \operatorname{Ad}_{\exp\Omega} W + (2X\Omega_z)^V\right) dz$$
$$(V^{\perp}): \quad \kappa' \cdot X = \left(2\Omega_z + (2X\Omega_z)^{\perp} + \lambda X \operatorname{Ad}_{\exp\Omega} W\right) dz - D'X.$$

Substitution implies

(33) 
$$\lambda[\operatorname{Ad}_{\exp\Omega} W, X]dz = 2\Big(\Omega_z + (X\Omega_z)^{\perp} - (X\Omega_z)^V X\Big)dz - D'X.$$

Conversely if (33) holds then  $\kappa' = (\lambda \operatorname{Ad}_{\exp \Omega} W + (2X\Omega_z)^V) dz$  and so  $\kappa'$  takes values in V. Comparing the  $\lambda^j$  coefficients on both sides of (33) we can solve for X inductively over j by at each stage requiring  $X_j \in \operatorname{im} \operatorname{ad}_{\operatorname{Ad}_{\exp \Omega} W}$  and

$$[\operatorname{Ad}_{\exp\Omega} W, X_1] = 2\Omega_z$$
$$[\operatorname{Ad}_{\exp\Omega} W, X_{j-1}]dz = 2\Big((X_j\Omega_z)^{\perp} - \sum_{s+l=j} (X_s\Omega_z)^V X_l\Big)dz - DX_k.$$

Define  $\nabla_{\lambda} = d + \operatorname{ad}_{\varphi_{\lambda}}$  and note that (7) says precisely that  $\nabla_{\lambda}$  is a flat connection in the trivial bundle  $\mathbb{C}/\Lambda \times \mathfrak{g}^{\mathbb{C}}$ . With X as above we have  $\nabla'_{\lambda}Y = 0$ . We wish to show that  $\nabla''_{\lambda}Y = 0$  also, as this will imply that Y satisfies the Killing field equation (13).

Define B by

(34) 
$$\nabla_{\lambda}''Y = (1+X)^{-1}B(1+X).$$

Using  $\operatorname{Ad}_{\exp\Omega} W = (1+X)Y(1+X)^{-1}$  and

$$\nabla_{\lambda}^{\prime\prime} \operatorname{Ad}_{\exp\Omega} W = \left[-\Omega_{\bar{z}} + \lambda^{-1} \operatorname{Ad}_{\exp-\Omega} \overline{W}, \operatorname{Ad}_{\exp\Omega} W\right] d\bar{z}$$

we obtain

$$Bd\bar{z} = [(-\Omega_{\bar{z}} + \lambda^{-1} \operatorname{Ad}_{\exp -\Omega} \overline{W}) d\bar{z} - \nabla_{\lambda}'' X (1+X)^{-1}, \operatorname{Ad}_{\exp \Omega} W]$$

which shows that B takes values in  $V^{\perp}$ .

As  $\nabla_{\lambda}$  is a flat connection we have commutativity of covariant derivatives and hence  $\nabla'_{\lambda} \nabla''_{\lambda} Y = 0$  which we write as

(35) 
$$-\nabla'_{\lambda}X(1+X)^{-1}B + \nabla'_{\lambda}B + B\nabla'_{\lambda}X(1+X)^{-1} = 0.$$

Since  $\nabla'_{\lambda}B = D'B + [2\Omega_z + \lambda \operatorname{Ad}_{\exp\Omega} W, B]dz$ , we can rewrite (35) as

(36) 
$$D'B = [\nabla'_{\lambda}X(1+X)^{-1} + (\lambda \operatorname{Ad}_{\exp\Omega} W - 2\Omega_z)dz, B].$$

From its defining equation (34) we know that B is of the form  $\sum_{j \leq d} \lambda^j B_j$ . We will show that B = 0. Suppose not, then there is some non-zero top coefficient  $B_d$ . Since X has only negative powers of  $\lambda$ , the  $\lambda^{d+1}$  term in (36) is

$$[\operatorname{Ad}_{\exp\Omega} W, B_d].$$

However we know that  $B_d \in V^{\perp}$  and hence it must be zero. Thus  $\nabla_{\lambda}'' Y = 0$  and Y satisfies the Killing field equation. From (32) we see that Y is of the form  $\sum_{j \leq 0} \lambda^j Y_j$  and furthermore  $Y_0 = \operatorname{Ad}_{\exp \Omega} W$ .

We now need to project this Y onto  $\Omega^{\sigma}(\mathfrak{g}^{\mathbb{C}})$  to get a solution to the Killing field equation in the correct loop algebra.

Representations of simple Lie algebras are completely reducible and we have identified  $\mathfrak{g}^{\mathbb{C}}$  with a subalgebra of  $\mathfrak{gl}(m,\mathbb{C})$  so it must have a complementary subspace in  $\mathfrak{gl}(m,\mathbb{C})$  which is invariant under the adjoint action of  $\mathfrak{g}^{\mathbb{C}}$ . This means there exists a projection map  $\pi : \Omega(\mathfrak{gl}(m,\mathbb{C})) \to \Omega(\mathfrak{g}^{\mathbb{C}})$  such that

$$d\pi(Y) = \pi(dY) = \pi([Y,\varphi_{\lambda}]) = [\pi(Y),\varphi_{\lambda}].$$

Thus we have that  $\pi(Y) \in \Omega(\mathfrak{g}^{\mathbb{C}})$  satisfies the Killing field equation. Furthermore  $Y_0 = \pi(\operatorname{Ad}_{\exp\Omega} W) = \operatorname{Ad}_{\exp\Omega} W$ . Set  $\tilde{Y} = \lambda Y = \sum_{j \leq 1} \lambda^j Y_{j-1}$  and note that  $\tilde{Y}_1 = Y_0 = \operatorname{Ad}_{\exp\Omega} W$ .

We want to project  $\tilde{Y}$  onto  $\Omega^{\sigma}(\mathfrak{g}^{\mathbb{C}})$ . Consider the map

$$\pi_j^{\sigma} := \frac{1}{k} (\operatorname{Id} + \epsilon^{-j} \sigma^j + \epsilon^{-2j} \sigma^{2j} + \ldots + \epsilon^{-(k-1)j} \sigma^{(k-1)j})$$

where  $\epsilon$  is the k-th primitive root of unity. This map  $\pi_j^{\sigma}$  projects any element in  $\mathfrak{g}^{\mathbb{C}}$  to its part in  $\mathfrak{g}_j$ . Thus we can define  $\pi^{\sigma} : \Omega(\mathfrak{g}^{\mathbb{C}}) \to \Omega^{\sigma}(\mathfrak{g}^{\mathbb{C}})$  by

$$\pi^{\sigma}(\sum_{j}\lambda^{j}\xi_{j}) = \sum_{j}\lambda^{j}\pi_{j}^{\sigma}(\xi_{j}).$$

Then  $\tilde{\xi} = \pi^{\sigma}(\tilde{Y})$  satisfies

$$d\tilde{\xi} = [\tilde{\xi}, \Omega_z + \lambda \operatorname{Ad}_{\exp\Omega} W)dz + (-\Omega_{\bar{z}} + \lambda^{-1} \operatorname{Ad}_{\exp-\Omega} \overline{W})dz]$$

and  $\tilde{\xi}_1 = \tilde{Y}_1 = \operatorname{Ad}_{\exp \Omega} W$ .

Now we may apply Lemma 5.1 to  $\tilde{\xi}$  to conclude the existence of a (real) polynomial Killing field  $\xi$  whose top term,  $\xi_d$ , is  $\operatorname{Ad}_{\exp\Omega} W$ .

The d-1 coefficient of  $\xi_z = [\xi, \Omega_z + \lambda \operatorname{Ad}_{\exp \Omega} W]$  is

$$(\operatorname{Ad}_{\operatorname{exp}\Omega} W)_z = [\operatorname{Ad}_{\operatorname{exp}\Omega} W, \Omega_z] + [\xi_{d-1}, \operatorname{Ad}_{\operatorname{exp}\Omega} W]$$

which implies

$$[\xi_{d-1} - 2\Omega_z, \operatorname{Ad}_{\exp\Omega} W] = 0.$$

Since W is a cyclic element and  $\xi_{d-1} - 2\Omega_z \in \mathfrak{t}$  we conclude  $\xi_{d-1} - 2\Omega_z = 0$  and hence  $\xi$  satisfies the theorem.

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